Behind the Mpemba paradox

Chang Qing Sun*
Nanyang Technological University; Singapore

Mpemba paradox results from hydrogen-bond anomalous relaxation. Heating stretches the O:H nonbond and shortens the H—O bond via Coulomb coupling; cooling reverses this process to emit heat at a rate depending on its initial storage. Skin ultra-low mass density raises the thermal diffusivity and favors outward heat flow from the liquid. As the source and central part of all lives, hydrogen bond (O:H—O, the ‘:’ is the lone pair and the ‘—’ bonding pair of electrons) plays key role in DNA and protein folding, drug-target binding, ion-channel activating/deactivating, gene delivering, cancel cell curing, messaging, and signaling, etc. However, response of hydrogen bond to temperature and coordination environment is crucial to these activities though mechanisms remain yet unknown. This article shows how the O:H—O bond relaxes in resolving the Mpemba paradox based on our extensive research.

Figure 1 illustrates (a) the potential paths and (b) the initial temperature dependence of the liner velocity of the O:H—O bond during thermal relaxation. The hydrogen bond consists of the intermolecular O:H nonbond (left-hand side of the H\(^+\) reference origin) and the intramolecular H—O bond (right-hand side) other than either of them alone. The O:H—O bond acts as an asymmetric, H\(^+\) bridged, and coupled oscillator pair with memory, whose cooperative relaxation stems anomalies of water ice such as ice floating, regelation, ice slippery, hydrophobic and tough water skin, Hofmeister effect, etc. An interplay of the O:H nonbond (van der Waals-like) interaction, the H—O covalent bond interaction, the O—O Coulomb repulsion and externally applied stimulus always dislocate O atoms in the same direction but by different amounts along the respective potential paths. The softer O:H nonbond relaxes more than the stiffer H—O does.

Generally, heating stores energy in a substance by stretching and softening all bonds involved. However, heating stores energy in water by stretching the O:H nonbond and simultaneously compressing the H—O bond (red line linked spheres in Figure 1a are in the hot state) by the joint interactions. Cooling reverses (blue line linked spheres) this process, analogous to suddenly releasing a coupled, highly deformed bungee pair, one of which is stretched and the other compressed. This process emits energy at a rate that depends on the deformation history of the bungee pair (i.e., how much they were stretched or compressed). The O:H—O bond exhibits memory with thermal momentum during cooling. The linear velocity of the H—O bond \(\Delta v/\Delta t\) in Figure 1b is the product of slopes for the known temperature dependence of H—O length \(d_\text{H}(\theta)\) and the measured initial temperature and time dependence of water temperature \(\theta(\theta_i, t)\) shown in Figure 1c as inset. Because of the Coulomb coupling of the O:H and H—O, the velocity \(\Delta v/\Delta t\) and energy rate \(\Delta E/\Delta t\) (\(x = L\) for O:H and \(x = H\) for H—O) can be readily derived but here we show the representative only for simplicity. Figure 1b indicates that the initially shorter H—O bond at higher temperature remains highly active compared to its behavior otherwise when they meet on the way to freezing.

Conversely, molecular undercoordination (with fewer than 4 neighbor as in the bulk) has the same effect of heating on O:H—O bond relaxation. H—O bond contracts and O:H expands, which shrinks molecular size and enlarges their separations with an association of polarization. Water skin forms such a supersolid phase that is elastic, hydrophobic, ice like, and less dense \((0.75\ \text{gcm}^{-3})\). The lower mass density raises the thermal diffusivity and favors outward heat flow from the liquid.

Figure 1c and d show the theoretical duplication of measured (insets) Mpemba paradox. The Mpemba paradox is characterized by: (i) hot water freezes faster than cold water under the same conditions; (ii) the liquid temperature \(\theta\) drops exponentially with cooling time \(t\); (iii) water skin is warmer than sites inside the liquid and the skin of hotter water is even warmer. Besides the thermal diffusivity and convection velocity involved in the non-linear Fourier thermal fluid transport equation, systematic examination of all possible boundary conditions of a 10-mm long, one-dimensional tube cell of water with a one-millimeter thick skin cooling from \(\theta_i\) to \(\theta_f = 0\ ^\circ\text{C}\) revealed the following:

1. Characterized by the crossing temperature in Figure 1c, the Mpemba effect happens only in the presence of the supersolid skin.
2. The Mpemba effect is sensitive to the source volume, the extent of skin supersolidity, skin radiation and the drain temperature \(\theta_i\).
3. The thermal convection has little effect on observations.

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*Correspondence to: Chang Qing Sun; Email: Ecqsun@ntu.edu.sg
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It is necessary to emphasize that the Mpemba effect happens only in circumstances where the water temperature drops sharply from $u_i$ to $u_f$ at the source–drain interface. Larger liquid volume may prevent this effect by heat-dissipation hindering. Any spatial temperature decay between the source and the drain could prevent the Mpemba effect. Conducting experiments under identical conditions is necessary to minimize artifacts such as radiation, source/drain volume ratio, exposure area, container material, etc. Therefore, conditions for the Mpemba effect are indeed very critical, which explains why it is not frequently observed.

Therefore, O:H—O bond memory and water-skin supersolidity determines the heat ‘emission-conduction-dissipation’ dynamics of the Mpemba paradox in the ‘source-path-drain’ cycle system and the skin-drain interface must be highly nonadiabatic.

Disclosure of Potential Conflicts of Interest

No potential conflicts of interest were disclosed.

References